

NLO EW Corrections for W^+Z scattering at the LHC

$pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj + X @ \mathcal{O}(\alpha^6)$ and $\mathcal{O}(\alpha^7)$ for $\sqrt{s} = 13$ TeV

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with:

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Vector-boson scattering in a nutshell

→ Scattering of two (massive) vector-bosons, e.g.:

- $W^\pm W^\pm \rightarrow W^\pm W^\pm$ (“like-sign W scattering”)
- $W^\pm Z \rightarrow W^\pm Z$

$$W^+ Z \rightarrow W^+ Z$$

$$\mathcal{M}_{\text{VBS}} =$$

$$pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj + X$$

$$\mathcal{M}_{\text{VBS@LHC}} =$$

Vector-boson scattering (VBS) physics program:

- Constrain anomalous quartic gauge couplings (with triple-gauge boson prod.)
 - Measure Higgs-vector-vector couplings, complementary to on-shell Higgs decay measurements
 - Probe EW symmetry breaking: interplay between triple and quartic gauge couplings and the Higgs boson(s); large cancellations for longitudinal VBS: ensures **tree-level unitarity**
- Precise prediction of the SM cross section needed

Longitudinal VBS: Tree-level (non-)unitarity

 $W^+Z \rightarrow W^+Z$
 $@ M_H = 125 \text{ GeV}$

$$\mathcal{M} = \underbrace{\text{diagram 1}}_{\mathcal{M}_4} + \underbrace{\text{diagram 2}}_{\mathcal{M}_s} + \underbrace{\text{diagram 3}}_{\mathcal{M}_u} + \underbrace{\text{diagram 4}}_{\mathcal{M}_H}$$

$$\mathcal{M}_4 \propto -s^2 - u^2 - 4su + 2(M_W^2 + M_Z^2) \frac{s^2 + 6su + u^2}{s + u} + \dots$$

$$\mathcal{M}_s \propto s^2 + 2su - 2M_W^2 \frac{3su + u^2}{s + u} - 2M_Z^2 \frac{2u^2 + 3su - s^2}{s + u} - \frac{M_Z^4}{M_W^2} s + \dots$$

$$\mathcal{M}_u \propto u^2 + 2su - 2M_W^2 \frac{3su + s^2}{s + u} - 2M_Z^2 \frac{2s^2 + 3su - u^2}{s + u} - \frac{M_Z^4}{M_W^2} u + \dots$$

$$\mathcal{M}_H \propto -\frac{M_Z^4}{M_W^2} \frac{t^2(t - 4M_W^2)(t - 4M_Z^2)}{(t - M_H)(t - 2M_W^2)(t - 2M_Z^2)} = -\frac{M_Z^4}{M_W^2} t + \dots$$

Assuming
 $|t| \gg M_H^2, M_Z^2, M_W^2$

$$\mathcal{M} = \mathcal{M}_4 + \mathcal{M}_s + \mathcal{M}_u + \mathcal{M}_H \propto 0 + \dots$$

Example from [Schwartz]

Longitudinal VBS: Tree-level (non-)unitarity

$$W^+Z \rightarrow W^+Z$$

$$@ M_H = \infty$$

$$\mathcal{M} = \underbrace{\text{diagram 1}}_{\mathcal{M}_4} + \underbrace{\text{diagram 2}}_{\mathcal{M}_s} + \underbrace{\text{diagram 3}}_{\mathcal{M}_u} + \underbrace{\text{diagram 4}}_{\mathcal{M}_H}$$

$$\mathcal{M}_4 \propto -s^2 - u^2 - 4su + 2(M_W^2 + M_Z^2) \frac{s^2 + 6su + u^2}{s + u} + \dots$$

$$\mathcal{M}_s \propto s^2 + 2su - 2M_W^2 \frac{3su + u^2}{s + u} - 2M_Z^2 \frac{2u^2 + 3su - s^2}{s + u} - \frac{M_Z^4}{M_W^2} s + \dots$$

$$\mathcal{M}_u \propto u^2 + 2su - 2M_W^2 \frac{3su + s^2}{s + u} - 2M_Z^2 \frac{2s^2 + 3su - u^2}{s + u} - \frac{M_Z^4}{M_W^2} u + \dots$$

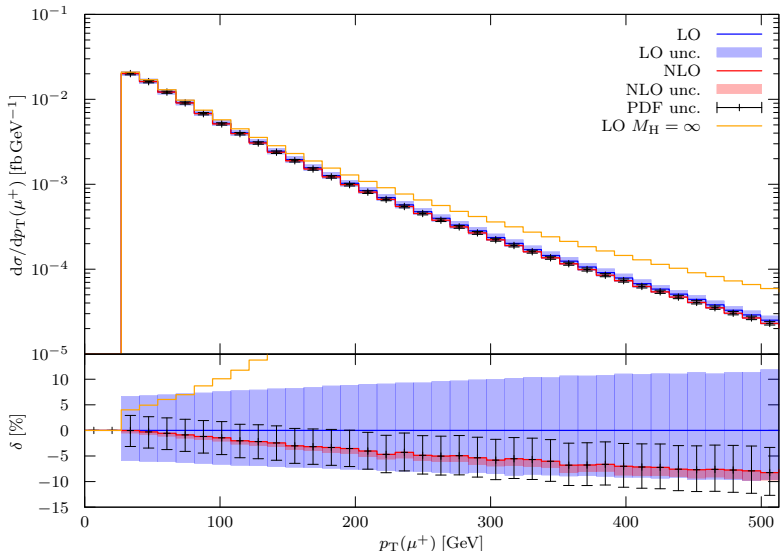
$$\mathcal{M}_H \propto -\frac{M_Z^4}{M_W^2} \frac{t^2(t - 4M_W^2)(t - 4M_Z^2)}{(t - M_H)(t - 2M_W^2)(t - 2M_Z^2)} = 0$$

$$\mathcal{M} = \mathcal{M}_4 + \mathcal{M}_s + \mathcal{M}_u + \mathcal{M}_H \propto -\frac{M_Z^4}{M_W^2} (s + u) + \dots$$

Example from [Schwartz]

→ $M_H = \infty$ estimates the maximal effect of different Higgs couplings

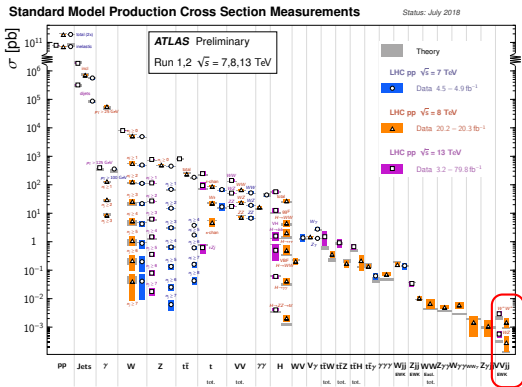
$$M_H = \infty \text{ for } pp \rightarrow e^+ \nu_e \mu^+ \nu_\mu jj + X$$



→ Large positive correction, range for extended Higgs sector

Experimental status for VBS, EW $pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj + X$

→ VBS processes are $\mathcal{O}(1 \text{ fb})$, need large \sqrt{s} and \mathcal{L} :



(dinosaur plot from the [\[ATLAS Collaboration\]](#))

- ATLAS 8 TeV: [\[CERN-EP-2016-017\]](#)
- ATLAS 13 TeV: Observ. with 5.6 σ sig. ($\mathcal{L} = 36.1 \text{ fb}^{-1}$) [\[ATLAS-CONF-2018-033\]](#)
- CMS 13 TeV: Meas. with 1.9 σ sig. ($\mathcal{L} = 35.9 \text{ fb}^{-1}$) [\[CMS-PAS-SMP-18-001\]](#)
- Easiest VBS channel is $W^+W^+ \rightarrow W^+W^+$, full NLO corrections available [\[Biedermann, Denner, Pellen\]](#)
- Next channel: $W^+Z \rightarrow W^+Z$

Discriminating between EW and QCD production

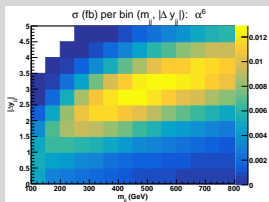
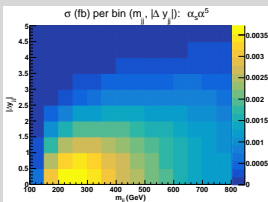
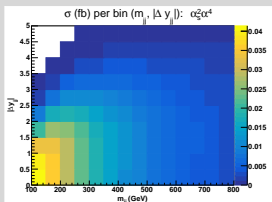
→ At LO, three different coupling orders:

$$\mathcal{M}_{\text{EW}} = \text{[diagram: EW production]} + \dots \quad \mathcal{M}_{\text{QCD}} = \text{[diagram: QCD production]} + \dots \rightarrow$$

- $\mathcal{O}(\alpha^4 \alpha_s^2)$: $|\mathcal{M}_{\text{QCD}}|^2$
- $\mathcal{O}(\alpha^5 \alpha_s^1)$: $2\Re\{\mathcal{M}_{\text{QCD}}^* \mathcal{M}_{\text{EW}}\}$
- $\mathcal{O}(\alpha^6 \alpha_s^0)$: $|\mathcal{M}_{\text{EW}}|^2$

How to maximize the EW production (“signal”)?

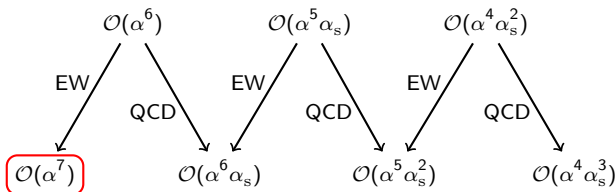
(pp → e⁺ν_eμ⁺ν_μjj + X)



(plots from VBSCAN WG1 report [Ballestrero, et. al.]

- Observables $M_{j_1 j_2}$ and $\Delta y_{j_1 j_2}$ are used to discriminate the QCD from the EW production
- In the fiducial PS region, EW > QCD for like-sign scattering (no initial-state gluons), for W⁺ Z-scattering: QCD > EW

LO and NLO for $pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj + X$



- All LOs presented in Sec. V.3 of the SM Les Houches 2017 report [Bendavid et. al.]: QCD (~80%) dominates over EW
 - Approx. $\mathcal{O}(\alpha^6 \alpha_s^1)$: [Bozzi, Jäger, Oleari, Zeppenfeld]
 - $\mathcal{O}(\alpha^4 \alpha_s^3)$ calculation available [Campanario, Kerner, Ninh, Zeppenfeld]
- $\mathcal{O}(\alpha^7)$ EW corrections desirable, because like-sign case shows large corrections (-16%)

Validation and checks

We performed **two independent calculations** for both $\mathcal{O}(\alpha^6)$ and $\mathcal{O}(\alpha^7)$:

“Freiburg”

- MEs from OpenLoops [Cascioli, Maierhöfer, Pozzorini]
- Loops evaluated with DD (COLI fallback) from COLLIER [Denner, Dittmaier, Hofer]
- General purpose Monte Carlo [CS]
- Dipole subtraction [Catani, Seymour] to regularize IR singularities
- PDFs from LHAPDF 6 [Buckley, et. al.]

“Würzburg”

- MEs from RECOLA [Actis, Denner, Hofer, Scharf, Uccirati]
- Loops evaluated with COLI (and DD) from COLLIER
- MoCaNLO [Feger] used by M. Pellen
- CS dipole subtraction with α -dependent dipoles [Nagy]
- PDFs from LHAPDF 6

Extensive checks:

- NLO virtuals checked against each other for 1000 PS points passing the cuts
- Integrated cross sections
- Each bin of 23 differential distributions, ca. 7800 bins

Fiducial phase space volume for $pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj + X$

Cuts chosen similar to the ATLAS

8 TeV-analysis [CERN-EP-2016-017]:

- At least two $R = 0.4$ anti- k_t jets with $p_T > 30$ GeV, $|\eta| < 4.5$, and $\Delta R_{j\ell} > 0.3$
- $M_{j_1 j_2} > 500$ GeV, **no $\Delta\eta_{j_1 j_2}$ cut**¹
- $p_{T,\ell} > 20$ GeV and $|y_\ell| < 2.5$
- $p_{T,\text{miss}} > 30$ GeV
- $|M_{\mu\bar{\mu}} - M_Z| < 10$ GeV
- $\Delta R_{\ell\ell} > 0.3$

Other:

- Photons recombined with charged particles using anti- k_t algorithm with $R = 0.1$
- PDFs: NNPDF30_nlo_as_0118_qed
- $\sqrt{s} = 13$ TeV

Complex mass scheme [Denner, Dittmaier, Roth, Wackerroth][Denner, Dittmaier, Roth, Wieders], input parameters:

- $G_\mu = 1.663\,787 \times 10^{-5} \text{ GeV}^{-2}$
- $M_W = 80.357\,97 \text{ GeV}$, $\Gamma_W = 2.084\,30 \text{ GeV}$
- $M_Z = 91.153\,48 \text{ GeV}$, $\Gamma_Z = 2.494\,27 \text{ GeV}$
- $M_H = 125.0 \text{ GeV}$, $\Gamma_H = 4.07 \times 10^{-3} \text{ GeV}$

with coupling calculated as:

$$\alpha = \frac{\sqrt{2}}{\pi} G_\mu M_W^2 \left(1 - \frac{M_W^2}{M_Z^2} \right)$$

Scale choice: $\mu_F = (1/2, 1, 2) \cdot M_W$

→ **No dependence on μ_R** , since processes do not depend on α_s !

¹Unused in the ATLAS 8 TeV-analysis, but used both in the ATLAS and CMS 13 TeV analyses

Feynman diagrams and partonic channels

$$\mathcal{M}_{\text{LO}} = \text{[Diagram 1]} + \text{[Diagram 2]} + \text{[Diagram 3]} + \text{[Diagram 4]} + \dots$$

Borns and Born-like:

- $uu \rightarrow e^+ \nu_e \mu^+ \mu^- du$ ($\sim 47\%$ XS),
- $du \rightarrow e^+ \nu_e \mu^+ \mu^- dd$ ($\sim 18\%$ XS),
- ... + 38 more: **bottleneck are virtuals**

Reals:

- $uu \rightarrow e^+ \nu_e \mu^+ \mu^- du \gamma$,
- $du \rightarrow e^+ \nu_e \mu^+ \mu^- dd \gamma$,
- ... + 38 more,

Not included, small or negligible:

- $\gamma\gamma \rightarrow e^+ \nu_e \mu^+ \mu^- (d\bar{u}/s\bar{c})$, and
- $bu \rightarrow e^+ \nu_e \mu^+ \mu^- db$,
- ... + 7 more, with resonant top-quarks:

$$\mathcal{M}_{\text{b-quarks}} = \text{[Diagram 5]} + \dots$$

Reals not yet calculated, expected to be small:

- $\gamma u \rightarrow e^+ \nu_e \mu^+ \mu^- du \bar{u}$,
- ... + 27 more,
- $\gamma\gamma \rightarrow e^+ \nu_e \mu^+ \mu^- d\bar{u} \gamma$, and
- $\gamma\gamma \rightarrow e^+ \nu_e \mu^+ \mu^- d\bar{u} \gamma$.

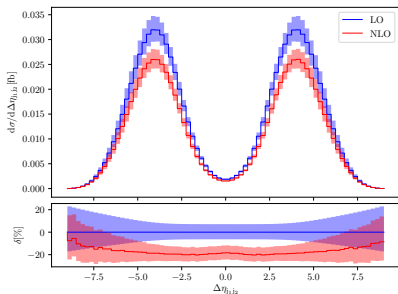
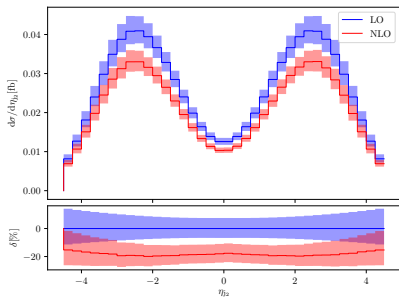
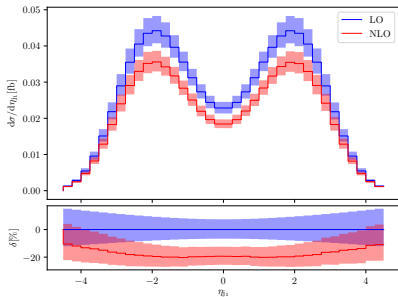
Integrated cross section

Integrated xs for $pp \rightarrow e^+ \nu_e \mu^+ \mu^- jj + X$ @ $\sqrt{s} = 13$ TeV for the fiducial PS volume:

LO [fb]	NLO [fb]	$\delta = \frac{\mathcal{O}(\alpha^7)}{\mathcal{O}(\alpha^6)}$ [%]
$0.2362^{+9.433\%}_{-8.022\%}$	$0.1899^{+8.356\%}_{-7.575\%}$	-19.6%

- Uncertainty is the range given by varying $\mu_F = M_W$ by (1/2,2)
- No dep. on $\alpha_s \rightarrow$ no dep. on μ_R
- **Huge corrections** (4–5 \times larger than e.g. EW corr. for di-boson prod.) on the integrated cross section, larger than even like-sign W-scattering (-16%)
- Corrections are even larger in specific regions of p_T distributions

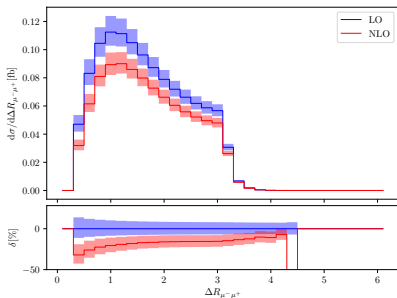
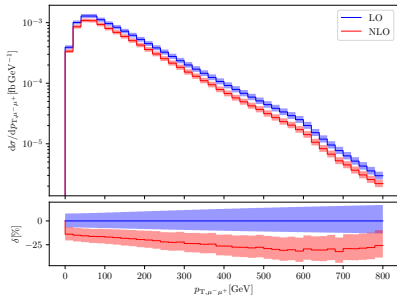
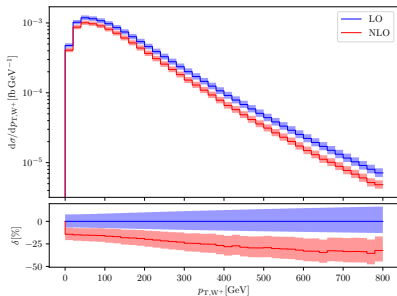
Jet pseudo-rapidities



- EW corrections flat in most PS regions
- Bands: large dependence on μ_F
- Peak at $\Delta\eta_{j_1 j_2} \approx 0$ suppressed because of $M_{j_1 j_2} > 500$ GeV cut:

$$\cosh \Delta\eta_{j_1 j_2} \approx \frac{M_{j_1 j_2}}{2p_{T,j_1} \cdot p_{T,j_2}} + \cos \Delta\phi_{j_1 j_2}$$

Leptonic observables



$$\Delta R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$$

- In p_{T,W^+} Sudakov logs further increase EW corrections
- $\Delta R_{\mu\bar{\mu}}$ is limited from above because $M_{\mu\bar{\mu}} \approx M_Z$ cut limits $\Delta\eta_{\mu\bar{\mu}} < 3.3$
- Kink at $p_{T,\mu\bar{\mu}} \approx \frac{2M_Z}{\Delta R_{\ell\ell}}$ caused by $\Delta R_{\ell\ell}$ cut

Summary

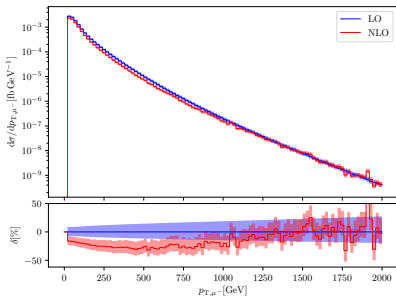
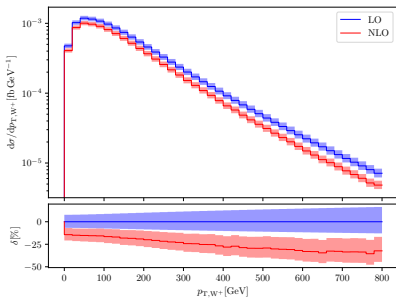
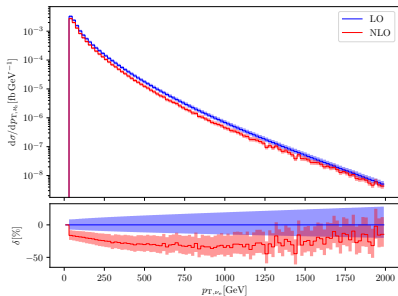
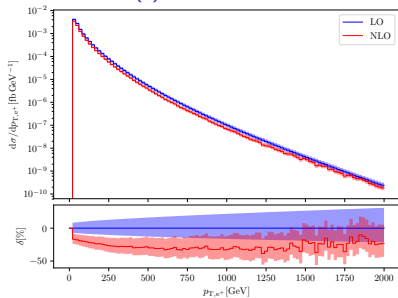
- Vector-boson scattering (VBS) important to constrain the actual Higgs-sector, complementary to Higgs-couplings measurements
- After W^+W^+ , W^+Z scattering is the next important channel for VBS
- Experiments are either measuring or already observing it
- EW corrections are available now, huge correction on the integrated cross-section: -20%

Acknowledgments

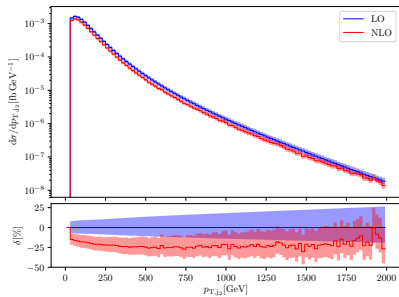
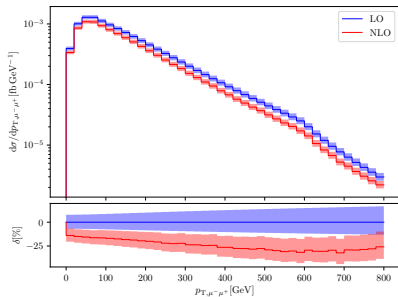
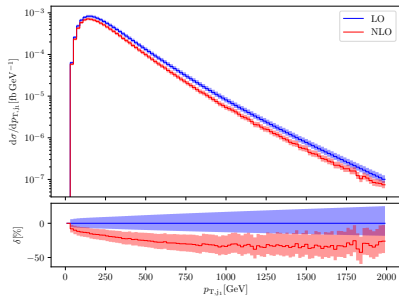
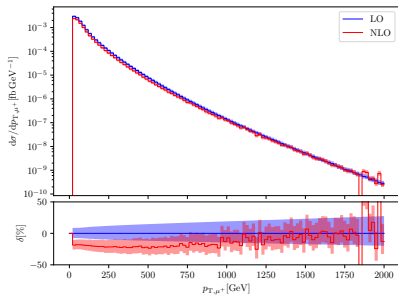
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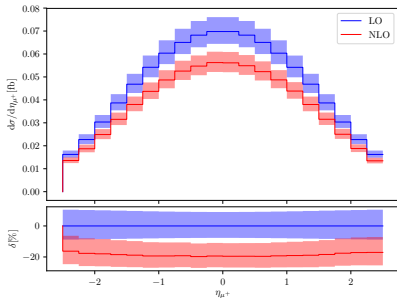
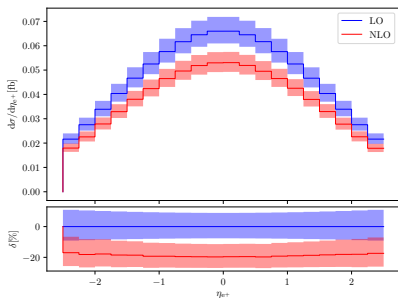
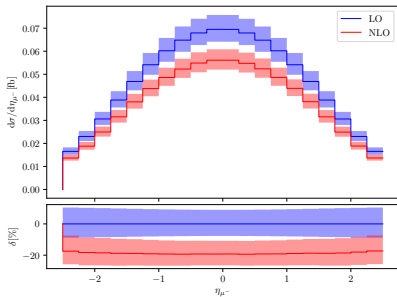
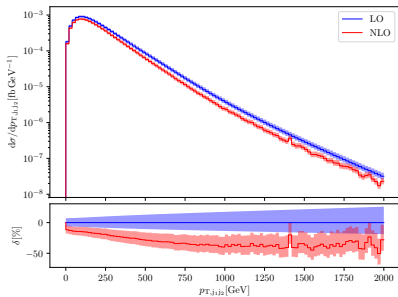
Distributions (I)



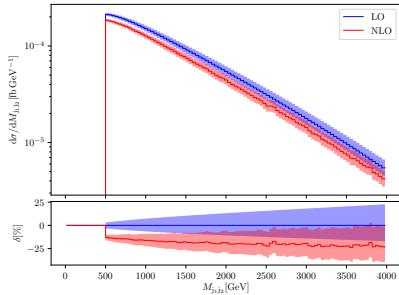
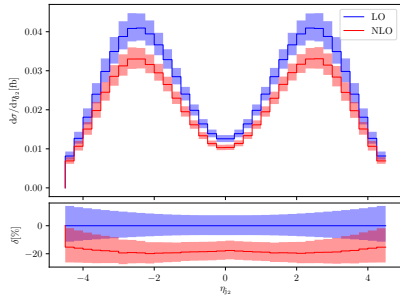
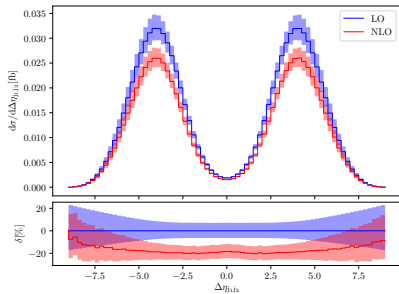
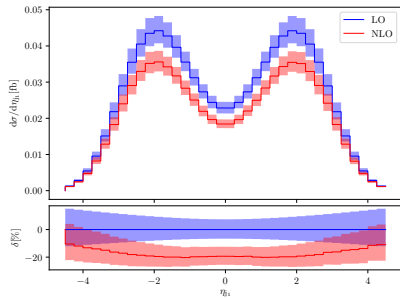
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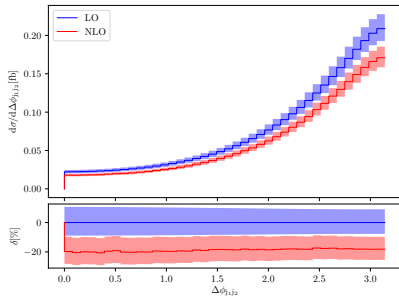
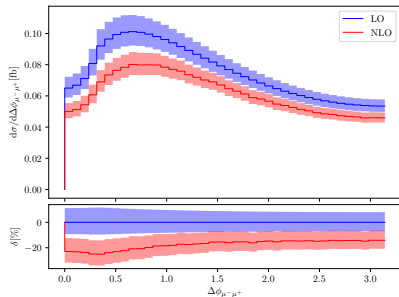
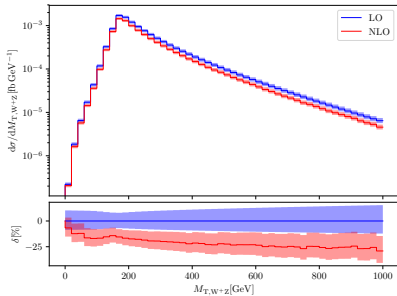
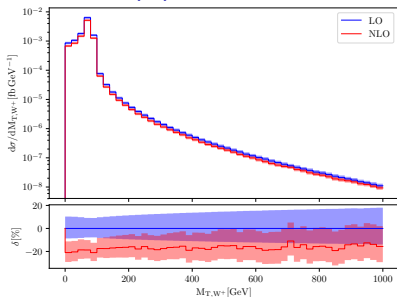
Distributions (III)



Distributions (IV)



Distributions (V)



Distributions (VI)

